



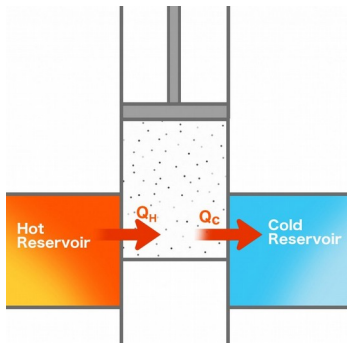
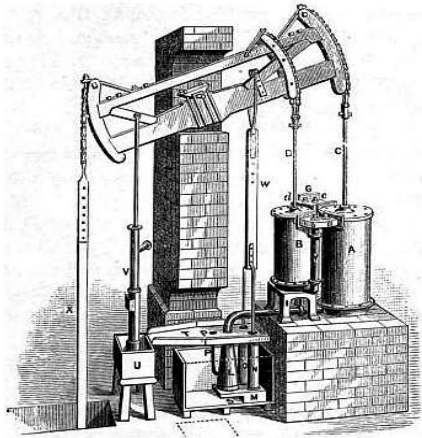
Thermodynamics of Open Chemical Reaction Networks

Massimiliano Esposito

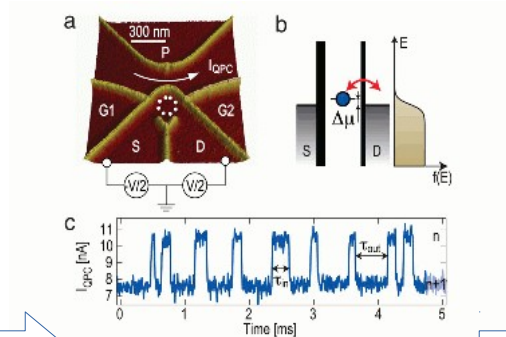
Riccardo Rao (PhD),
Gianmaria Falasco (postdoc), Matteo Polettini (researcher)

Harvard, November 13, 2019

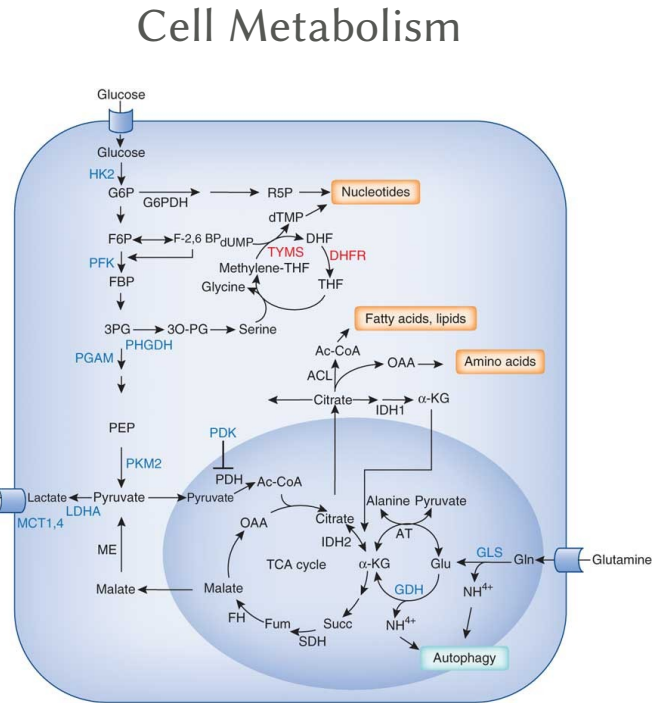
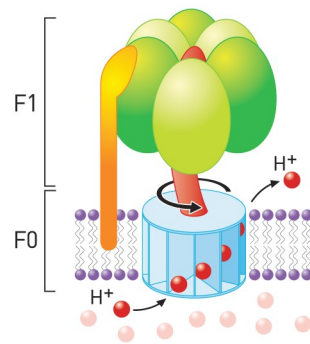
Motivation



Classical thermodynamics



Stochastic thermodynamics



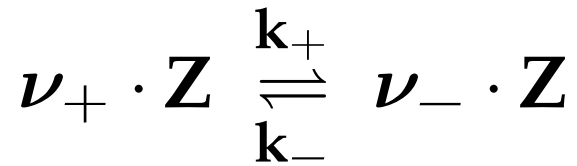
Thermodynamics of Chemical Reaction Networks

“[Thermodynamics] is the only physical theory of universal content which I am convinced will never be overthrown, within the framework of applicability of its basic concepts.” A. Einstein

Outline

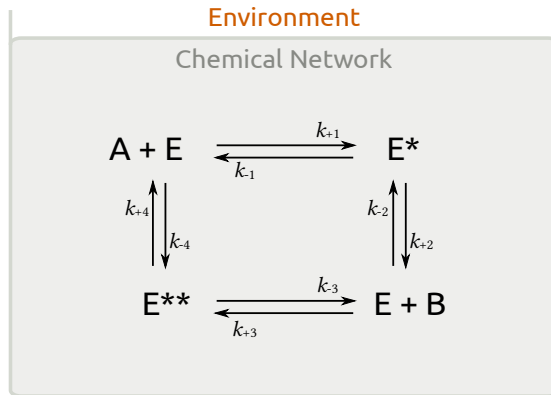
- Deterministic Dynamics of Open CNRs
- Thermodynamics of Open CRNs
 - First and Second Law
 - Role of Topology
- Reaction-Diffusion
- Concluding Remarks

Dynamics of Closed CRNs



Stoichiometric matrix

$$\nu_- - \nu_+ =: \mathcal{S}$$



$$\mathcal{S} = \begin{pmatrix}
 1 & 2 & 3 & 4 \\
 -1 & 1 & -1 & 1 \\
 1 & -1 & 0 & 0 \\
 0 & 0 & 1 & -1 \\
 -1 & 0 & 0 & 1 \\
 0 & 1 & -1 & 0
 \end{pmatrix} \begin{array}{l} \text{E} \\ \text{E}^* \\ \text{E}^{**} \\ \text{A} \\ \text{B} \end{array}$$

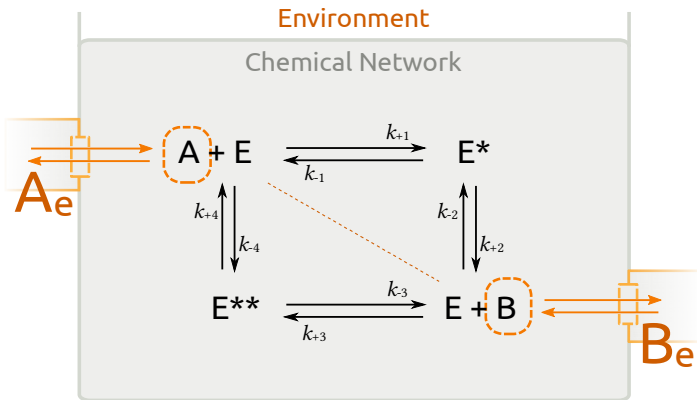
Rate equations $\mathbf{d}_t[\mathbf{Z}] = \mathcal{S}\mathbf{J}$

Reactions

Ideal Dilute Solution + Elementary Reactions
= Mass action kinetics

$$\left\{ \begin{array}{l} \mathbf{J} = \mathbf{J}_+ - \mathbf{J}_- \\ \mathbf{J}_\pm = \mathbf{k}_\pm[\mathbf{Z}] \cdot \nu_\pm \end{array} \right.$$

Dynamics of Open CRNs



$$\nu_+ \cdot \mathbf{Z} \xrightleftharpoons[k_-]{k_+} \nu_- \cdot \mathbf{Z}$$

Stoichiometric matrix

$$\nu_- - \nu_+ =: \mathbf{S} = \begin{pmatrix} \mathbf{S}^X \\ \mathbf{S}^Y \end{pmatrix}$$

$$[\mathbf{Z}] = \begin{pmatrix} [\mathbf{X}] \\ [\mathbf{Y}] \end{pmatrix} \begin{array}{l} \text{Internal} \\ \text{Chemostatted} \end{array}$$

Rate equations

$$d_t[\mathbf{X}] = \mathbf{S}^X \mathbf{J}$$

$$d_t[\mathbf{Y}] = \mathbf{S}^Y \mathbf{J} + \mathbf{I}$$

Reactions Exchange

$$\mathbf{S} = \begin{pmatrix} -1 & 1 & -1 & 1 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & -1 \\ -1 & 0 & 0 & 1 \\ 0 & 1 & -1 & 0 \end{pmatrix} \begin{array}{l} \mathbf{S}^X \\ \mathbf{S}^Y \end{array}$$

Autonomous CRNs $d_t[\mathbf{Y}] = 0$

Thermodynamics of CRNs

Enthalpy of formation

Entropy of formation

Ideal Dilute Solutions
Local equilibrium

$$\mu = \underbrace{h^\circ}_{\text{Enthalpy of formation}} - \underbrace{T s^\circ}_{\text{Entropy of formation}} + RT \ln[\mathbf{Z}]$$

μ° Standard-state chemical potential

Local Detailed Balance

$$\ln \frac{k_+}{k_-} = -\frac{\mu^\circ \cdot S}{RT}$$

Elementary Reactions

“0th law of Thermodynamics”: Closed CRN relax to equilibrium

$$\mathbf{J}^{\text{eq}} = \mathbf{J}_+^{\text{eq}} - \mathbf{J}_-^{\text{eq}} = 0 \quad \text{detailed-balance}$$

First and Second Law

CRN Enthalpy:

$$H = h^\circ \cdot [Z]$$

CRN Entropy:

$$S = (s^\circ - R \ln[Z]) \cdot [Z] + \overset{\text{total concentration}}{R[Z]}$$

1st law
Enthalpy Balance

$$d_t H = h^\circ \cdot \mathcal{S}J + h_Y \cdot I = h^\circ \cdot \mathcal{S}J + \underbrace{Ts_Y \cdot I + \mu_Y \cdot I}_{\text{Heat Flow } \dot{Q}} + \underbrace{\mu_Y \cdot I}_{\text{Chemical Work } \dot{W}_{\text{chem}}}$$

2nd law
Entropy Balance

$$\dot{\Sigma} = d_t S - \underbrace{\dot{Q}/T}_{\text{Entropy change in (thermal \& chemical) reservoirs}}$$

Entropy production:
(total entropy change)


$$T\dot{\Sigma} = -\mu \cdot \mathcal{S}J = (J_+ - J_-) \cdot RT \ln \frac{J_+}{J_-} \geq 0$$

Work principle: Equilibrium of Closed CRN

Non-Eq. Gibbs free energy $G = H - TS$

$$T\dot{\Sigma} = \dot{W}_{\text{chem}} - d_t G \geq 0$$

$$G = G_{\text{eq}} + RT \underbrace{\mathcal{L}([Z] | [Z]_{\text{eq}})}_{\text{“Relative entropy”}}$$

Equilibrium of closed CRN 

$$= [Z] \cdot \ln \frac{[Z]}{[Z]_{\text{eq}}} - ([Z] - [Z]_{\text{eq}}) \geq 0$$

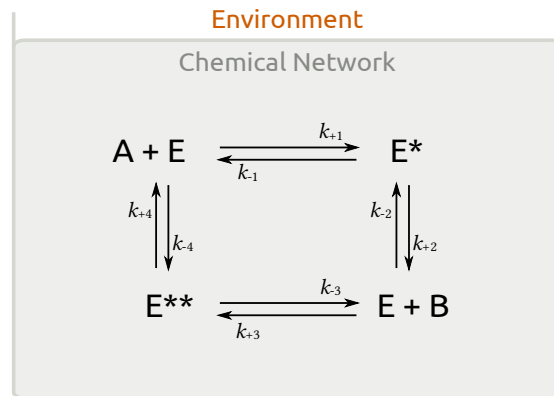
In a closed CRN, $W_{\text{chem}} = 0$, G is minimized by the dynamics until the CRN reaches equilibrium

Conservation Laws & Cycles in Closed CRNs

Conservation Laws: $\ell_\lambda \cdot \mathcal{S} = 0$

Cycles: $\mathcal{S} \mathbf{c}_\alpha = 0$

$$\mathbf{L}_\lambda = \ell_\lambda \cdot [\mathbf{Z}]$$



$$\mathcal{S} = \begin{pmatrix} 1 & 2 & 3 & 4 \\ -1 & 1 & -1 & 1 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & -1 \\ -1 & 0 & 0 & 1 \\ 0 & 1 & -1 & 0 \end{pmatrix} \begin{matrix} \text{E} \\ \text{E}^* \\ \text{E}^{**} \\ \text{A} \\ \text{B} \end{matrix}$$

$$\ell_1 = (1 \quad 1 \quad 1 \quad 0 \quad 0)$$

$$\mathbf{L}_1 = [\text{E}] + [\text{E}^*] + [\text{E}^{**}]$$

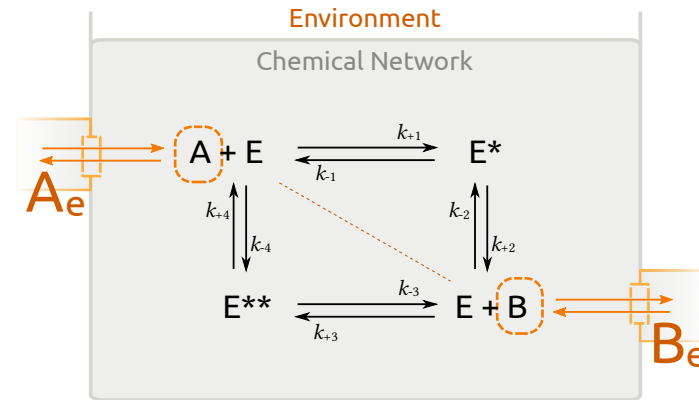
$$\ell_2 = (0 \quad 1 \quad 1 \quad 1 \quad 1)$$

$$\mathbf{L}_2 = [\text{E}^*] + [\text{E}^{**}] + [\text{A}] + [\text{B}]$$

$$\mathbf{c} = (1 \quad 1 \quad 1 \quad 1)^T$$

Conservation Laws & Cycles in Open CRNs

$$S = \begin{pmatrix} -1 & 1 & -1 & 1 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & -1 \\ -1 & 0 & 0 & 1 \\ 0 & 1 & -1 & 0 \end{pmatrix} \begin{matrix} S^X \\ \\ \\ S^Y \end{matrix}$$



Opening may break some conservation laws

$$\lambda \rightarrow \lambda_b, \lambda_u \quad \text{unbroken: } \ell_1 = (1 \quad 1 \quad 1 \quad 0 \quad 0) \quad L_1 = [E] + [E^*] + [E^{**}]$$

$$\ell_{\lambda_b} \cdot S^X \neq 0 \quad \text{broken: } \cancel{\ell_2 = (0 \quad 1 \quad 1 \quad 1 \quad 1)} \quad L_2 = [E^*] + [E^{**}] + [A] + [B]$$

Opening may create emergent cycles

$$\begin{cases} S^X \mathbf{c}_e = 0 \\ S^Y \mathbf{c}_e \neq 0 \end{cases} \quad \mathbf{c} = (1 \quad 1 \quad 1 \quad 1)^T$$

$$\text{emergent: } \mathbf{c}_e = (1 \quad 1 \quad 0 \quad 0)^T$$

$$\#Y = \# \lambda_b + \# \mathbf{c}_e$$

$$Y \rightarrow Y_b, Y_f$$

$$1 = 1 + 0$$

$$2 = 1 + 1$$

$$2 = 2 + 0$$

Entropy Production shaped by Topology

$$\# = \#Y - \#\lambda_b$$

Fundamental Forces Emergent Cycle Affinities

$$\mu_{Y_f} - \mathbb{L}\mu_{Y_b}$$

$$-\mu_Y \cdot S^Y c_\epsilon$$

$$\mathcal{F}_{Y_f} \cdot \mathbf{I}_{Y_f} = \mathcal{A}_\epsilon \cdot \mathcal{J}_\epsilon$$

$$T\dot{\Sigma} = \dot{W}_d + \dot{W}_{nc} - d_t \mathcal{G} \geq 0$$

Nonconservative
Work

Driving Work

$$\dot{W}_d = -\sum_{\lambda_b} [\partial_t f_{\lambda_b}(\mu_{Y_b})] L_{\lambda_b}$$

NonEq semigrand Gibbs free energy

$$\mathcal{G} = G - \sum_{\lambda_b} f_{\lambda_b}(\mu_{Y_b}) L_{\lambda_b}$$

$$\mathcal{T}\dot{\Sigma} = \dot{W}_d + \dot{W}_{nc} - d_t \mathcal{G} \geq 0$$

Some special class of CRNs

Detailed Balanced CRNs: $\mathcal{F}_{Y_f}, \mathcal{A}_\epsilon = 0 \rightarrow \dot{W}_{nc} = 0$

DB & autonomous: $\mathcal{T}\dot{\Sigma} = -d_t \mathcal{G} \geq 0$ Relaxation to Eq.: \mathcal{G} minimized by dynamics

Non-Eq. Steady State: $\mathcal{T}\dot{\Sigma} = \mathcal{F}_{Y_f} \cdot \mathbf{I}_{Y_f} = \mathcal{A}_\epsilon \cdot \mathcal{J}_\epsilon \geq 0$ $\#Y - \#\lambda_b$ Force-Flux pairs
(autonomous)

Work principle

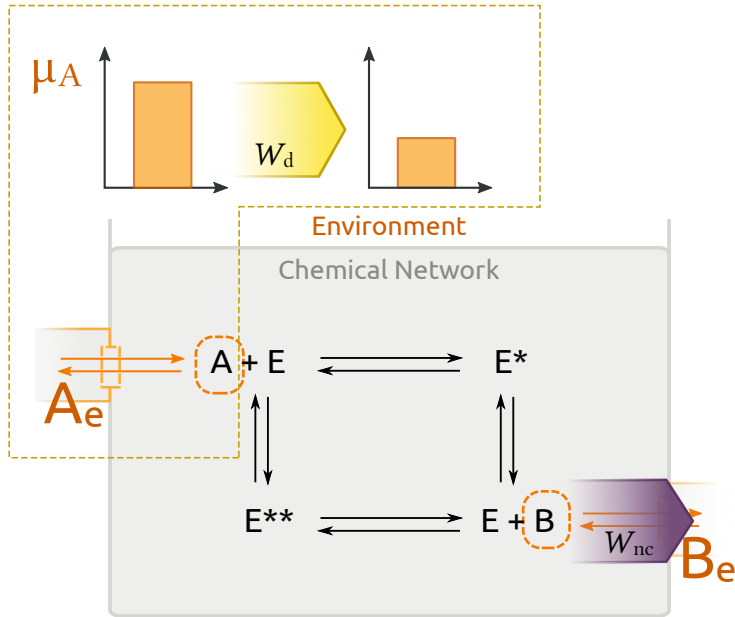
$$\mathcal{G} = \mathcal{G}_{eq} + RT \mathcal{L}([\mathbf{X}] | [\mathbf{X}]_{eq}) \geq 0$$



Equilibrium of open CRN

$$\dot{W}_d + \dot{W}_{nc} \geq \Delta \mathcal{G}_{eq} + \mathcal{L}([\mathbf{X}] | [\mathbf{X}]_{eq}) \rightarrow \text{Minimal work to generate a Non-Eq distribution}$$

Example



$$L_b = [E^*] + [E^{**}] + [A] + [B]$$

	Driving Work	Nonconservative Work	Noneq. semigrand Gibbs free energy
	$-\underbrace{[\partial_t \mu_A] L_b}_{\text{red}}$	$\underbrace{(\mu_B - \mu_A) I_B}_{\text{blue}}$	$\underbrace{\mathcal{G} = G - \mu_A L_b}_{\text{red}}$
$T\dot{\Sigma}$	$= \dot{W}_d$	$+ \dot{W}_{nc}$	$- d_t \mathcal{G} \geq 0$

Adding Diffusion

Reaction–Diffusion Equations

$$d_t[\mathbf{X}]_r = -\nabla \cdot \mathbf{J}_r^{\mathbf{X}} + \mathbf{S}^{\mathbf{X}} \mathbf{j}_r$$

$$d_t[\mathbf{Y}]_r = -\nabla \cdot \mathbf{J}_r^{\mathbf{Y}} + \mathbf{S}^{\mathbf{Y}} \mathbf{j}_r + \mathbf{I}_r$$

Diffusion
Reactions
Exchange

Mass-action kinetics

$$\mathbf{j}_r^\pm = \mathbf{k}_\pm [\mathbf{Z}]_r^{\circ \nu_\pm}$$

Diffusion: Fick's Law

$$\mathbf{J}_r = -\mathbb{D} \nabla [\mathbf{Z}]_r$$

↑
Diffusion coefficients

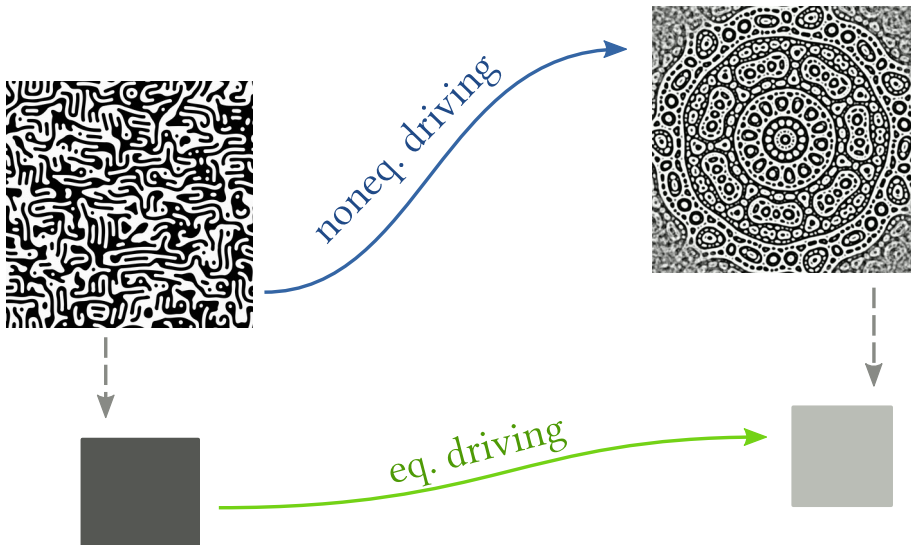
$$W_d + W_{nc} = \Delta \mathcal{G}_{eq} + RT \Delta \mathcal{L} + T \Sigma$$

$$\geq 0$$

$$= \int_{\mathcal{V}} d\mathbf{r} \left[[\mathbf{Z}]_r \circ \ln \frac{[\mathbf{Z}]_r}{[\mathbf{Z}]_{eq}} - ([\mathbf{Z}]_r - [\mathbf{Z}]_{eq}) \right] \geq 0$$

Spatial structuring takes work:

$$\mathcal{L}([\mathbf{Z}]_{patt} | [\mathbf{Z}]_{eq}) \geq \mathcal{L}([\bar{\mathbf{Z}}] | [\mathbf{Z}]_{eq})$$



Final remarks

The big issue of coarse graining

Non-Eq. thermodynamics needs to describe **all** the degrees of freedom that are out-of-equilibrium must be described. We need coarse graining strategies which preserves thermodynamic consistency.

Ex: Coarse-graining Biocatalysts [[Wachtel, Rao & Esposito, *New J. Phys.* **20** 042002 \(2018\)](#)]

Useful?

- Performance of energy storage and energy conversion (from molecular motors to metabolic networks)

Ex: Driven synthesis [[Penocchio, Rao & Esposito, *Nature Communications*, **10** 3865 \(2019\)](#)]

- Cost of Chemical Information Processing

Ex: Turing patterns [[Falasco, Rao & Esposito, *Phys. Rev. Lett.* **121**, 108301 \(2018\)](#)]

Chemical waves [[Avanzini, Falasco & Esposito, arXiv:1904.08874](#)]

Stochasticity?

Stochastic Thermodynamics of CRNs (nonlinear) [[Rao & Esposito, *J. Chem. Phys.* **149**, 245101 \(2018\)](#)]

Thermo at stochastic vs deterministic level: yes at steady state for complex balanced CRNs

[[Polettini, Wachtel & Esposito, *J. Chem. Phys.* **143**, 184103 \(2015\)](#)]

Phase transitions [[Lazarescu, Cossetto, Falasco & Esposito, *J. Chem. Phys.* **151**, 064117 \(2019\)](#)]